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TITLE APPLICATION OF MUON SPIN RELAXATION EQUIPMENT TO THE VORTEX DYNAMICS OF HIGH TEMPERATURE SUPERCONDUCTORS

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We discuss the potential use of muon spin relaxation (μ SR) technique to investigate the dynamics of the magnetic flux lines (vortices) in the mixed state superconductors. We argue that with the high critical temperature and large amsotropy found in some high temperature superconductors, one can use μ SR to observe both static and dynamic behaviors of vortices. Recent experiments on Bi₂Si₂CaCu₂O₈ are interpreted in this context.

In high temperature superconductors, it is well known that magnetic flux lines (vortices) are mobile at temperatures close to the critical temperature (T) due to thermal fluctuations? However, the nature of the low remperature vortex state, remains to be clarified. The difficulty lies mainly in the techniques employed in investigating the dynamics at low temperature. All the previous methods relied on measuring the response of vortices under external perturbation or non equilibrium configuration," and it is unclear if any of these methods can really be used to study the low temperature regimes where vortices are not very mobile and responsive. One exception is the Bitter pattern decoration experiment However, this method is limited to very low external field, and the measurement is done only on the surface of the sample of

The muon spin relaxation (μ SR) experiment, on the other hand, is a technique that is largely non-perturbative to the system, and explores the bulk. However, due to the relatively short lifetime of muons (about 2.2 μ s), it has not been a very effective tool in studying the vortex eastion. Recally, vortices look "frozen" within such time scales in ordinary superconductors. In fact, this has been used to advantage in measuring the penetration length λ of the scalaterials—as well as high T_c systems without much vortex motion.

In certain high I, materials (such as 16,81,0 a) (12,000), however, the version sees seem to be strongly mobile, clear

ly due to the higher anisotropy (they have almost layered structures), and μSR may be capable of investigating their dynamic behavior. Let us consider the necessary criterion for this to be possible

The most important quantity in a μ SR measurement is the second moment of the held distribution $\Delta B^{\perp} =$ $\ell(B - \langle B \rangle)^2$) where $\langle \cdots \rangle$ stands for the spatial and time average over the muons' lifetimes. In the case of static for near static evoltices, this average is same as just the spatial average. However, in the case of very mobile vortices, ΔB^2 , is reduced due to the time averaging process. One can see this by simply considering the peak (i.e. the highest) time averaged held in the system a muon can observe. When the vortices are frozen, it is simply the field at a vortex core, whereas when the vortices are mobile, this held distribution is smeared by the vortex motion a Note that mesons do not move with you tices to The same applies to the minimum held except that it is shifted upwards. Thus ΔB^2 must become smaller when vortices move.

A more claborate analysis has shown in fact that even the temperature dependence of ΔBC may be altered. depending on how much motion is allowed. In stead of a standard 1 Δ^4 type temperature dependence inegative curvature with temperatures at may relief to positive curvature.

For the vortex motion to manifest itself blocking starnecessary for a vortex to make one oblibe convibrates a significant fraction of the inter-vortex distances within a muon's lifetime. (The fraction needed is about 1.4 at $T_{\rm c}$, and the rms distance from the original position having $T^{1/2}$ like temperature dependence, consistent with either diffusive or vibrating motion.) We can roughly estimate from the resistivity measurement of Ref. 1 if this condition is met. Consider the typical speed at which the vortices for vortex discs, in the case of layered su perconductors) move in these experiments: Icm's at $10^{5} \sim 10^{6} \text{ A} \text{ m}^{2} \text{ at } T \approx 80 \text{ K for BijSrjCaCujOs}$. From this, and assuming that vortices are independent of each other and that Einstein's relation holds, we can estimate the diffusion constant for this material. We find it to be about $10^9 \, \text{Å}^2 \, \mu\text{s}$, a number too large to be reasonable. Clearly, a correction is necessary to account for the more microscopic theory of energy dissipation associated with vortex motion. However, the fact that this number turns out to be so large is an indication that the vortices are very mobile in Bi₂Sr₂CaCu₂O₈, and that this motion can be seen in µSR experiment.

The actual experiment on Bi₂Sr₂CaCu₂Os has recently shown. that $-\Delta B^2$, indeed shows positive curvature with temperature at H=3 and 4 kG. We have compared these data with our theory, and obtained a satisfactory agreement if we assume that the vertices move about 250 V within a muon lifetime at I_c . The overall scale of $-\Delta B^2$ is strongly influenced by the amount of vortex disorder present in the sample as it is cooled down to the lowest possible temperature. Therefore, the penetration length V is difficult to estimate for these materials. Nevertheless, we estimate this to be about $3000 \ \Delta^{2}$

Our study also showed that the third moment of the field distribution $\Delta B^{\prime\prime}$ is just as important in studying the low temperature vortex dynamics. It can be shown that the third moment which is a measure of the asymmetry of the field distribution, is small when you cas are mobile. This is simply due to the fact that no main can measure the core magnetic field, and that the field distribution noticinger has a long both magnetic field tail a sociated with the core. Such a tail is only observable when the vortex motion is small. Thus, by looking at ΔBC at different temperatures one can be cate the Treezing temperature of vortice. We find this to be about 30 K for 40, 80, our 40 K below which ΔBC

shows a rapid increase from ≈ 0 with lowering temperature. It seems to indicate that it is not a first order phase transition (ergo, not melting), and it is found to be consistent with a rapid crossover of the thermal activation picture as well as a second order or Kosterlitz-Thouless transition. If the thermal activation model applies, this "freezing" temperature must be dependent on the time and length scales involved in the experiment, (for μ SR, they are $2.2\,\mu$ s and a few hundred Λ).

Finally, since the field distribution may show a temperature dependent asymmetry as mentioned above, it is clearly not sufficient to simply fit a gaussian decay to a μSR signal. It, in fact, does not even give a good fit for Bi₂Sr₂CaCu₂O₈. We suspect that this will be the case for any mixed state superconductors with high critical temperature and large anisotropy, and we also believe that it is essential to check if the quality of the fit is satisfactory before (ΔB^2) etc. are computed. Otherwise, the result would be inaccurate and misleading.

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